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by

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Summary

The hydrodynamics code LASNEX is used to model the laser and ebeam generated plasma-plume experiments. The laser used has a wavelength of 1 μ m and the FWHM spot size is 1 mm. The total laser energy is 160 mJ. The simulation shows that the plume expands at a velocity of about 6 cm/ μ s. The e-beam generated from the Experimental Test Accelerator (ETA) has 5.5 MeV and FWHM spot size ranges from 2 to 3.3 mm. From the simulations, the plasma plume expansion velocity ranges from about 3 to 6 mm/ μ s and the velocity increases with decreasing spot size. All the simulation results reported here are in close agreement with experimental data.

Description of the LASNEX Simulations and Discussion of the Results

We have performed simulations to model the laser and e-beam generated plasma-plume experiments. In this memo, we present the simulation of the target expansion using the radiation hydrodynamics code LASNEX. The typical target in the experiment is made of tantalum with a density of $16.6 \, \text{g/cm}^3$ and a thickness of 1 mm. The laser used to generate the plasma plume has a wavelength of 1 μ m. The Gaussian pulse has a FWHM pulse length of 10 ns. The FWHM spot size is 1 mm and the total laser energy is 160 mJ. The electron beam from the ETA has kinetic energy of 5.5 MeV and a current of 1.7 kA. The pulse length is 70 ns. The rise and fall time of the pulse are 10 ns and the central flat-top portion has a 50 ns duration. The FWHM spot size ranges from 2 to 3.3 mm.

(I) Simulation of the plasma plume generated by laser illumination

The laser used to generate the plasma plume has a wavelength of 1 μ m. The pulse shape is Gaussian with FWHM pulse length of 10 ns and the FWHM spot size is 1 mm. The total laser energy is 160 mJ. The plot of the laser energy vs time is shown in Fig. 1.

Because the laser intensity is relatively weak and the skin depth is very shallow, the thickness of the outermost zone of the mesh used for this

simulation must be very thin, e.g. in the order of several nanometer. The configuration of the mesh is shown in Fig. 2a and the enlarged plot of the mesh near the vacuum-target interface is shown in Fig. 2b. In Fig. 2b, the thickness of the outermost zones of the mesh is too thin to be resolved in the figure.

The density contours are shown in Fig. 3a and the temperature distribution and the average ionization Z of the plasma plume at 20 ns are shown in Figs. 3b, respectively. Note that the electron-ion collision time τ_{ei} is substantially shorter than the laser pulse duration. Therefore, $T_i = T_e$.

The leading-edge expansion velocity vs time is plotted in Fig. 4. At about 100 ns, the leading edge expansion is about 90% of the terminal velocity of about 6 cm/ μ s, which is within the range of uncertainty of the experimentally measured velocity by Houck (Ref. 1).

Note that the leading edge velocity is less than 6 cm/ μ s if the velocity is obtained by dividing the distance expanded by the time in Fig. 3b. This is because the expansion velocity in early time is less than 6 cm/ μ s.

(II) Simulation of the plasma plume as a function of e-beam spot size

In this set of simulations, we have chosen three different e-beam spot radii --- 2.0, 2.5, and 3.2 mm. The energy deposited into the target by the e-beam is simulated by putting an energy source in the target. The initial configuration of the mesh is shown in Fig. 5. In order to speed up the simulation, we only simulate the left-half of the target. The electron beam is assumed to come in from the left.

The energy deposited vs time in the left-half of the slab is shown in Fig. 6a. The leading edge expansion velocity vs time, for the calculation with 2.0 mm spot radius, is shown in Fig. 6b. The temperature distribution and the Z contours of the plasma plume are shown in Figs. 7a and 7b, at 70 and 300 ns, respectively. At the end of the pulse at 70 ns, the maximum temperature inside the target is a little above 3 eV.

To measure the plume velocity, the forward Faraday cup for the plume velocity measurement is located at a distance of 25 cm from the target at an angle of 30° from the axis (Ref. 2). Therefore, the time it takes for the leading edge of the plume to reach the Faraday cup is in the order of tens of microseconds. This time is greater than the transit time (about 250 ns) for the leading edge of the plume to increase from zero to the maximum velocity as shown in Fig. 6b. Therefore, the velocity measured using the Faraday cup should be close to the maximum velocity shown in Fig. 6b multiplied by cos30°. The leading-edge expansion velocities, at an angle of 30° from the beam axis, obtained from the simulation, are

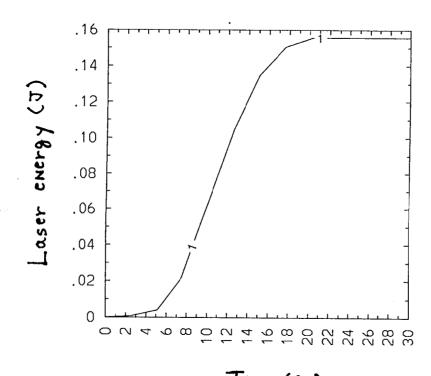
plotted vs the inverse of spot size in Fig. 8. The experimentally measured expansion velocities by Houck (Ref. 2) are also plotted here. The simulation result shows that the leading edge plasma velocity ranges from about 3.3 to 6.0 mm/ μ s as the spot radius decreases from 3.2 to 2.0 mm. This figure shows that the simulation and the experimental results are in good agreement.

Conclusion

The simulation results of the plasma plume velocities for the experiments described here are in good agreement with the experimentally measured data. This good agreement established the fact that LASNEX is a valuable tool for simulating laser and e-beam generated plasma plumes.

References:

- (1) T. Houck, personal communication, Lawrence Livermore National Laboratory, 1999.
- (2) T. Houck, et al., "Faraday Cup Measurements of the Plasma Plume Produced at an X-Ray Converter" LINAC98.



Time (ns)

Fig. 1

Mesh Configuration

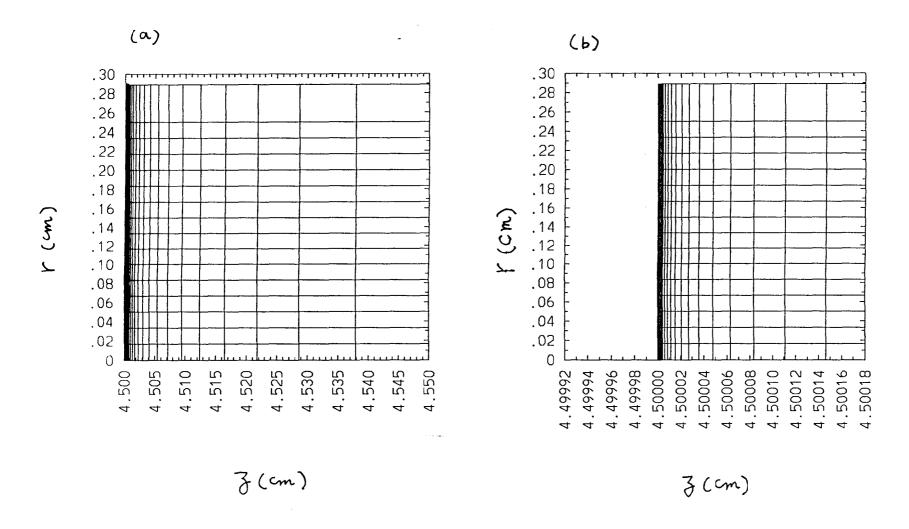


Fig. 2

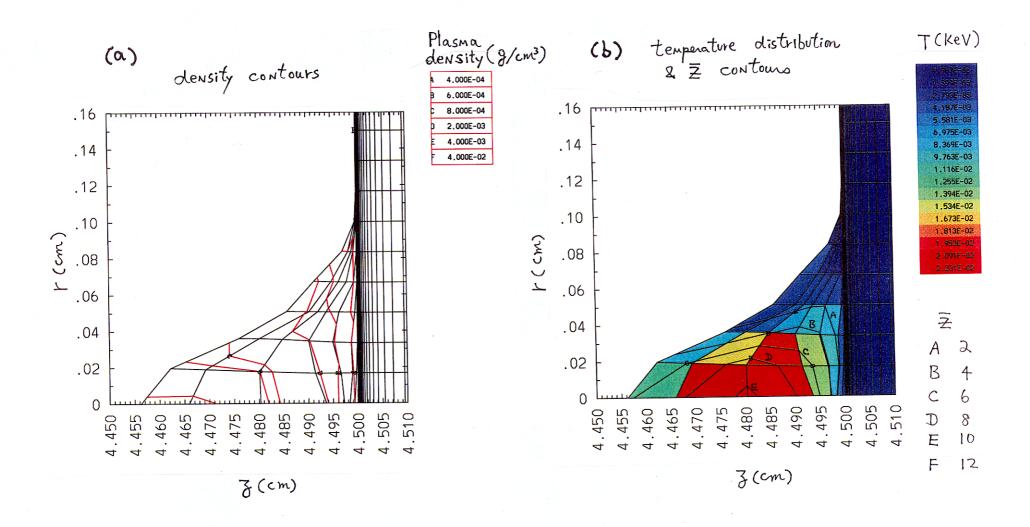


Fig. 3

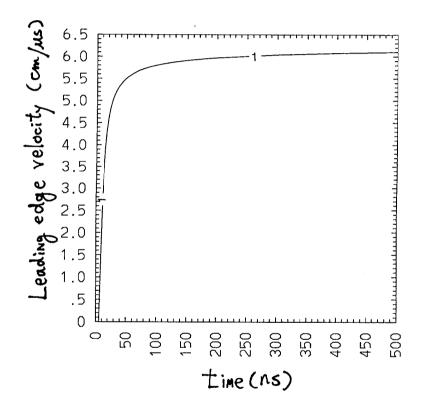


Fig. 4

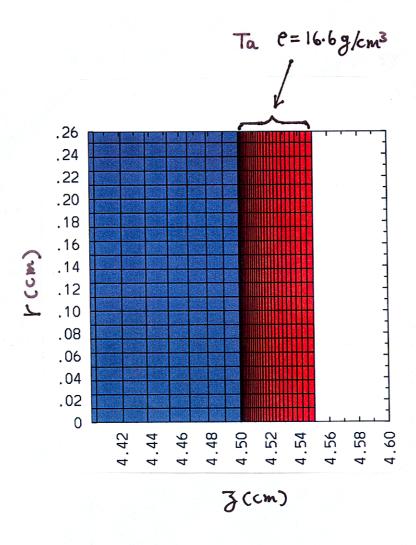


Fig. 5

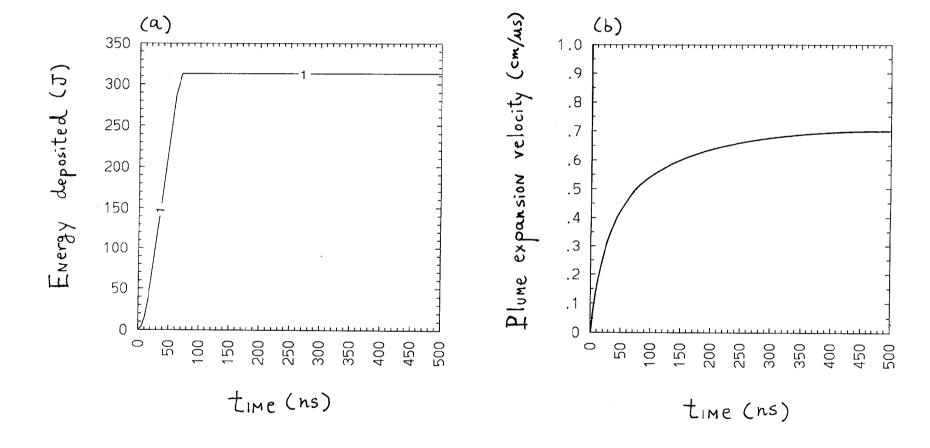


Fig. 6

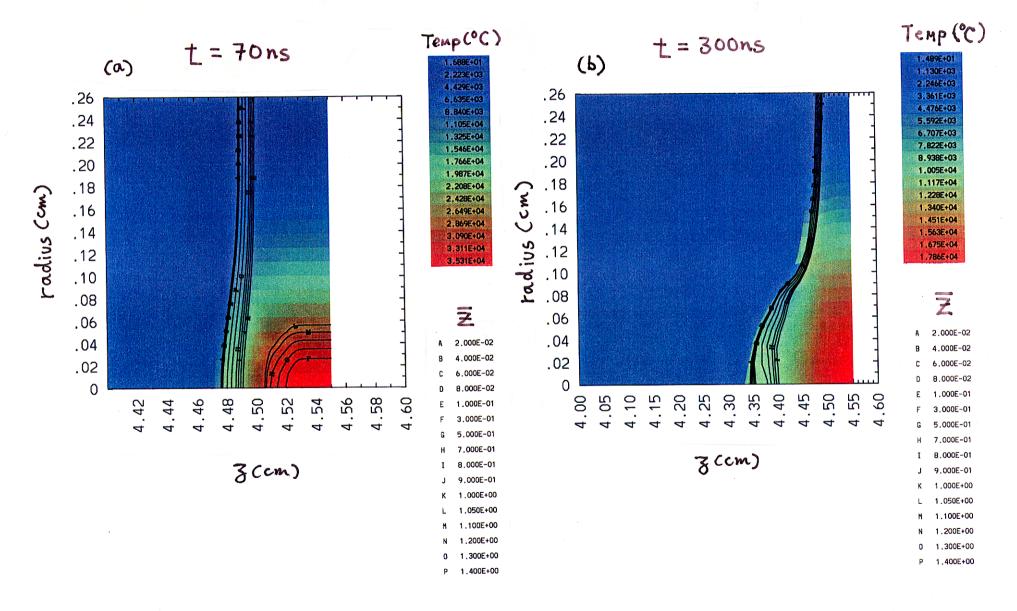


Fig. 7

Fig. 8